

**CORRIGENDUM TO "PHASE TRANSITIONS IN A PIECEWISE
EXPANDING COUPLED MAP LATTICE WITH LINEAR
NEAREST NEIGHBOUR COUPLING"**

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The proofs in our paper [1] need two corrections that we present below. The results remain valid as published. The first correction concerns the order in which some parameters in our construction have to be chosen; the second one fills a gap in the argument which extends the results from piecewise linear Markov maps to smooth circle maps.

1. CHOICE OF THE PARAMETER γ

The constant β appearing in **lemma 2.2.(a)**¹ has to be taken independently of γ , contrarily to what is written in the published version. The same has to be checked for the corresponding constant B in **lemma 4.1** and in **(5.3)**.

We need the independence of β and B from γ to perform the computation at the end of **section 5 (pp 2205–2206)**: since the constants K_1 and K_2 depend only on B , $K = 3(K_1 + \frac{1}{2}(1 + K_2)\beta)$ can in this case be chosen independently of γ and k . This is essential to conclude because the values of κ such that $\kappa K \frac{1}{2} \text{Var}(h_0) \leq \Delta^8$ depend (trivially) on K . Once the value of κ fixed, one has to choose γ small satisfying **lemma 2.2.(c)**, then k large satisfying **lemma 2.2.(a-b)**, without modifying K .

One can indeed fix β independently of γ thanks to the facts that $\tilde{\tau}$ is a piecewise affine Markov map and that γ does not appear in the lengths of images of its affinity intervals. Indeed, choose for any $\kappa > 0$ a k such that $\inf |(\tilde{\tau}^k)'| \geq \frac{2v}{\kappa}$. Since $\tilde{\tau}$ is piecewise affine the classical proof of Lasota-Yorke inequality (see for example proposition 2.1 in [2]) may be simplified to give

$$(1.1) \quad \text{Var}(P_{\tilde{\tau}}(f)) \leq v \text{Var}(P_{\tilde{\tau}^k}(f)) \leq \kappa \text{Var}(f) + \frac{2v}{\min_i |\tilde{\tau}^k(I_i)|} \int |f| dm,$$

where $\{I_i\}_i$ are the affinity intervals of $\tilde{\tau}^k$. But, as the affinity intervals of $\tilde{\tau}$ form a Markov partition for $\tilde{\tau}$, all intervals $\tilde{\tau}^k(I_i)$ occur as images $\tilde{\tau}(J)$ of affinity intervals J of $\tilde{\tau}$. Then an inspection of the specific form of the map $\tilde{\tau}$ shows that the constant

$$(1.2) \quad \beta = \frac{2v}{\min_i |\tilde{\tau}^k(I_i)|} = \frac{2(1-\eta)}{\delta}$$

depends on η and δ but is independent from γ and k . We can then conclude for the tensor product as in lemma 3.2 of [2], and for B as previously, since the modification of β that results in B will only be due to the coupling.

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¹All the boldfaced references are related to the published version of the paper.

2. THE CASE OF SMOOTH MODIFICATIONS

The formulation of **remark 2.3** is a bit misleading and the proof of **theorem 1 for the map $\bar{\tau}^3$** in **section 7** incomplete. We use in fact the specific construction of the local map in the proof of the exponential estimate from **section 5**, and it can not be directly transferred to a smooth modification of it.

More specifically, the inclusion **(5.12)** is obtained thanks to the calculations of **section 3**, in particular the two summarized observations at its end. These observations remain valid for $\tilde{\tau} = \frac{1}{v}\tilde{\tau}^k$, since $\tilde{\tau}$ has the same Markov partition structure as $\bar{\tau}$, but cannot be checked for a perturbation $\bar{\tau}$ such that

$$(2.1) \quad \rho := |||P_{\bar{\tau}} - P_{\tilde{\tau}}||| = \sup \left\{ \int |(P_{\bar{\tau}} - P_{\tilde{\tau}})(f)| dm : \text{Var}(f) \leq 1 \right\}$$

is small.

This difficulty can however be overcome, since formula **(5.12)** is only used in the "probabilistic" computation **(5.13)**. We explain below how this computation has to be modified when the local map $\bar{\tau}$ is a smooth modification of $\tilde{\tau}$. For all notations which are taken from the published version, we assume that the local map is $\bar{\tau}$. We denote also \bar{T} (resp. \tilde{T}) the 4-fold direct product of $\bar{\tau}$ (resp. $\tilde{\tau}$).

We want to compute

$$(2.2) \quad \int_{B_0(\mathbf{x}_{\Pi(i)})} h_0(y) dy = \int_{B_1(\mathbf{x}_{\Pi(i)})} P_{\bar{T}}^3 P_{\Phi_\epsilon^{\{t_1\}}} \left(1_{\bar{\Gamma}(i,t_1)} P_{\tilde{T}}^3 P_{\Phi_\epsilon^{\{t_1-1\}}} P_{\epsilon|\mathbf{x}_{\Pi(i)}}^{t_1-1} h_0 \right) (y) dy,$$

where

$$(2.3) \quad \bar{\Gamma}(i,t_1) := \{ \xi \in I^{U(i)} : (i, 0) \text{ is an error site for } \bar{T}^3 \circ \Phi_\epsilon^{\{t_1\}} \text{ at } (\xi, T_{\epsilon, \Pi(i)}^{t_1}(\mathbf{x}_{\Pi(i)})) \},$$

so that $E_{(i,t_1)} = T_{\epsilon, \mathbf{x}_{\Pi(i)}}^{-t_1}(\Gamma(i,t_1))$. We then need to replace:

- the six terms $P_{\bar{T}}$ by $P_{\tilde{T}}$, making an error of the order of $\rho K_3 \text{Var}(h_0)$, where K_3 is the sum of six uniform controls on norms of operators acting on BV spaces (since the operators $P_{\bar{T}}$ and $P_{\tilde{T}}$ satisfy a uniform Lasota-Yorke inequality).
- the term $1_{\bar{\Gamma}(i,t_1)}$ by its equivalent for \tilde{T} , $\check{\Gamma}(i,t_1) := \{ \xi \in I^{U(i)} : (i, 0) \text{ is an error site for } \tilde{T}^3 \circ \Phi_\epsilon^{\{t_1\}} \text{ at } (\xi, T_{\epsilon, \Pi(i)}^t(\mathbf{x}_{\Pi(i)})) \}$.

We notice that $\bar{\Gamma}(i,t_1)$ can be written as the disjoint union of 8 terms of the type $\bar{\Gamma}_{(i,t_1)}^{(1)} = \{ \xi : \xi_i \geq 0, \xi_{i+e_1} \geq 0, \xi_{i+e_2} \geq 0 \text{ and } (\bar{T}^3 \circ \Phi_\epsilon^{\{t_1\}}(\xi))_i \leq 0 \}$. One can then evaluate the difference (denoting $\check{\Gamma}_{(i,t_1)}^{(1)}$ the equivalent term for \tilde{T})

$$(2.4) \quad \begin{aligned} & \int |1_{\bar{\Gamma}_{(i,t_1)}^{(1)}} - 1_{\check{\Gamma}_{(i,t_1)}^{(1)}}| h(\xi) d\xi \\ & \leq \int 1_{\{\xi_i \leq 0\}} |(P_{\bar{T}}^3 - P_{\tilde{T}}^3) P_{\Phi_\epsilon^{\{t_1\}}} (1_{\{\xi_i \geq 0, \xi_{i+e_1} \geq 0, \xi_{i+e_2} \geq 0\}} h)|(\xi) d\xi \\ & \leq K_4 \rho \text{Var} h \end{aligned}$$

thanks to (2.1). The error made when replacing $1_{\bar{\Gamma}(i,t_1)}$ by $1_{\check{\Gamma}(i,t_1)}$ in (2.2) is then of the order of $\rho 8K_2K_4 \text{Var}(h_0)$.

We can now use the equivalent formulation of **(5.12)** for \check{T} ,

$$(2.5) \quad \check{\Gamma}_{(i,t_1)} \subseteq \{\xi : \xi \notin G\} \cup \{\xi : \check{T}\Phi_\xi^{\{t_1\}}(\xi) \notin G\} \cup \{\xi : \check{T}^2\Phi_\xi^{\{t_1\}}(\xi) \notin G\},$$

to compute the remaining integral

$$(2.6) \quad \int_{B_1(\mathbf{x}_{\Pi(i)})} P_{\check{T}}^3 P_{\Phi_\xi^{\{t_1\}}} \left(1_{\check{\Gamma}_{(i,t_1)}} P_{\check{T}}^3 P_{\Phi_\xi^{\{t_1-1\}}} P_{\epsilon|\mathbf{x}_{\Pi(i)}}^{t_1-1} h_0 \right) (y) dy$$

exactly as in **(5.13)**. One can finally conclude by choosing κ and ρ such that

$$(2.7) \quad (\kappa K + \rho(K_3 + 8K_2K_4)) \frac{1}{2} \text{Var}(h_0) \leq \Delta^8.$$

REFERENCES

- [1] J.-B. Bardet, G. Keller, *Phase transitions in a piecewise expanding coupled map lattice with linear nearest neighbour coupling*, *Nonlinearity*, **19** (2006), 2193–2210.
- [2] G. Keller, C. Liverani, *A spectral gap for a one-dimensional lattice of coupled piecewise expanding interval maps*, in: *Dynamics of Coupled Map Lattices and of Related Spatially Extended Systems* (Eds.: J.-R. Chazottes, B. Fernandez), *Lecture Notes in Physics* **671** (2005), pp. 115–151, Springer Verlag.

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